Superconducting Detector Magnets

Herman ten Kate

Content: 1. Concepts 2. Superconductors 3. Design of the CMS solenoid 4. The making of ATLAS 5. Future Collider Detectors



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How to discover new (elementary) particles?

 \checkmark Use E = mc² to produce particles from a package of energy.

We need E, an energy production unit (accelerator-collider), and an experiment to look at the shower of particles produced (detector).







Collision energy *E_{TeV} ≅ 0.3 B_T R_{km}* 9 T & 4.6 km → 14 TeV

Circular Collider:

Many magnets & few cavities, need higher magnetic field for a smaller ring High energy but growing synchrotron radiation losses ($\propto E^4/R$) High luminosity by a high bunch repetition rate Main bill is for the cryogenics for running the compressors to get 4 K.



Linear Collider:

Few magnets but nearly all cavities, need efficient RF power production A higher gradient will give a shorter machine Single shot, requiring a very small cross-section for high luminosity Main bill Is for the RF power.

Example: Large Hadron circular Collider

Exploring the energy frontier between up to 13-14 TeV using proton-proton & Pb-Pb collisions

CMS 🛄

LHC ring, 27 km circumference

ALICE

HE Physics and Superconductivity

LHC (and many other accelerators) can not be realized without extensive use of Superconductivity and High Quality Magnets



No Higgs without Superconductivity !

Large HEP detector magnets of the past...



Omega, medio 1972

BEBC, medio 1973

... and present detectors, CMS and ATLAS



CMS (2008)



ATLAS (2008)





Concept: why magnetic field in detectors

How to analyze the shower of particles ? We need:

- Track reconstruction
- Energy measurement (in calorimeters)
- Charge identification in magnetic field
- Momentum measurement in magnetic field.







Information yield:

- left turn => positively charged particle
- right turn => negative particle
- curvature => momentum



Concept: charged particle tracking

Example: tracking in the CMS Solenoid and iron return yoke



Concept: type of magnet used

- There are 3 principle magnet layouts for particle bending
- Choice depends on type of experiment and " 4π " or single direction fixed target, or even a combination of these, all variants exist.



Dipole magnet mainly vertical B Solenoid + yoke mainly axial B Toroid + Solenoid Tangential + axial B

Concept: sizing the detector

What determines the size of the generic " 4π " detector and magnetic field?

Radial thickness

- is the summation of:
- + tracking length inner detector
- + thickness of the solenoid
- + radial build of the calorimeters
- + tracking length
- + thickness of shielding iron yoke

Axial length

is the summation of:

<complex-block>

INNER TRACKER

CRYSTAL ECAL

HCAL

- + "catch angle" in forward directions sizing the length of the solenoid
- + thickness of iron shielding.

Concept: sizing the detector

What counts: momentum resolution! Particle with charge q and momentum p_t, travels through field B, is bent by Lorentz force:

 $\boldsymbol{F} = \boldsymbol{q} \left(\boldsymbol{E} + \boldsymbol{v} \boldsymbol{x} \boldsymbol{B} \right) \quad (\boldsymbol{E} \cong \boldsymbol{0})$

in the transverse direction, radius R, sagittal s:

$$s = \frac{L}{8R} = \frac{qBL^2}{8p_t}$$

and momentum resolution: $\frac{\partial n}{\partial t}$

$$p/p_t \propto p_t/0.3BL^2$$



p _t (GeV/c)	s [mm] @ B=1T, L=1m
1000	0.037
100	0.37
10	3.7
1	37

Keeping minimum the resolution for higher collision energies, so higher momenta, requires to scale the detector up with BL² !

- ! 10 times more energy \rightarrow 2xB and $\sqrt{5}$ =2.4x tracking length, say \approx diameter!
- ! and the axial length grows accordingly!

Thus: detectors grow in size with the colliding energy.

Concept: more requirements

- (1) Momentum resolution \rightarrow sufficient BL².
- (2) For physics we need B, not the magnet (!), though a rewarding challenge for magnet engineers!



 → Minimum thickness of coils to minimize particle scattering (especially when the calorimeters are put outside the central solenoid!)
 Material of choice: in general all AI, low density, inside the calorimeters.

- (3) Hermetically closed detector catching all particles.Minimum lost sphere for magnet services and supporting structures.
- (4) Full integration of magnets with detectors interleaved and supported.
- (5) Always working to avoid loss of data.

Requiring high operational margins in terms of temperature and current.

- (6) Unique and not replaceable (can not really be repaired).Very robust design with large margins and high level of redundancy.
- (7) And yes, low cost as well! NbTi at 4.5 K.



Pro's of Solenoids:

 B_{peak}/B_o ~ 1.2, so maximum B-yield for a given peak field, optimum use of superconductor.

- Cylindrical windings, easy and self-supporting.
- Forces, hoop stress and axial compressive stress are taken within the coil body, easy to optimize, symmetric and low heat in-leak.
- Windings can be supported by an outer support cylinder, also used as heat sink enabling conduction cooling of the coil.
- Coil can be thin and thus high transparency.
- Long track record of experience in scaling up.

Con's of Solenoids:

- Field not optimal for bending, not perpendicular to trajectories
- Massive iron flux return yoke, iron dominated, system very heavy.
- Less challenging....



Pro's of Toroids:

- Field perpendicular to trajectory, optimal bending and clean concept
- No iron yoke, so much lighter, but larger



- B=0 on the beam, no interference with beam and other parts.
- Challenging....

Con's of Toroids:

- B_{peak}/B_o ~4; given NbTi limits, only some 2 tesla can be used
- B=0 on the beam, thus toroids can not be used for inner detector
- Thus toroids can be used in combination with a central solenoid
- $B \propto 1/r$, so less uniform
- Forces not self-sustaining due to straight legs, need more stiffness
- Limited experience, but the largest detector magnet is a toroid!



2. Superconductors for detector magnets

Practical superconductors Basic properties Stability requirements Minimum Propagation Zone High Currents and Cables



From materials to magnets



How to make performing multi-kA conductors that guarantee the magnet not to quench or degrade ?

\rightarrow We need to understand and control the entire chain

- An under developed area of research, but essential to avoid surprises and degraded magnet performance
- Striking examples exist of missing understanding putting large projects at risk

Superconductors for magnets







Cubic alloy, isotropic





Tc: 11 K Bc₂: 13 T

Very well developed ~1 €/ kA m Barrel Toroid Conductor: 65 kA at 5 T

- 1.25 mm dia. NbTi/Cu strand, 2900 A/mm² at 5T
- 40 strands Rutherford cable, ~1700 A / strand
- Co-extruded with high purity AI (RRR>1500)
- Intermetallic bonding Cu-Al is required
- For the Barrel Toroid, size $57 \times 12 \text{ mm}^2$,
- 56 km made
- Production by 2 suppliers
- For the End Cap Toroids, size 41 x 12 mm²,
- 26 km made
- For the Central Solenoid, size 30 x 4.3 mm²
- 9 km made (Ni/Zn doped Al for higher Y-stress)







Coils and Superconducting Windings

As argued before, we need:

- 1 5 T, so we use NbTi
- thin and transparent, so we use AI
- simple cooling and robust mechanics.
 This caused an evolution of detector magnet design since some 40 yrs.

We see:

- Al stabilized Rutherford cables made from NbTi/Cu strands.
- 1-4 layer coils, often wound inside a supporting cylinder taking the hoop stress.
- Conduction cooled by thermo-siphon or forced He flow cooling at 4.5 K through Al tubes on the support cylinder.



Typical coil windings (ATLAS solenoid)



ATLAS Solenoid 2.5 T

Critical temperature, field dependency

Superconducting Phase (J_c vs. B and T).

For maintaining the superconducting state, the conductor must operate below the critical surface determined by critical current, magnetic field and temperature.

For NbTi the critical area is bounded by:

 $T_c(B=0) = 9.2 \text{ K} \text{ and } B_{c2}(T=0) = 14.5 \text{ T}$ $B_{c2}(T) = B_{c2}(0) [1 - (T/9.2)^{1.7}]$ $T_c(B) = Tc(0) [1 - (B/14.5)]^{0.59}$



 $B_{c2}(4.2 \text{ K}) = 10.7 \text{ T}$ $T_{c}(5 \text{ T}) = 7.16 \text{ K}$

Similar relations are found for Nb_3Sn and BSCCO 2212 and 2223.

Temperature margin, T_{cs}

When a transport current flows, the onset of resistance is is further reduced from T_c to T_{cs} , the current sharing temperature

 $T_{cs}(B,I) = T_b + (T_c(B) - T_b) (1 - I/I_c)$ $T_{cs}(5 T,I_c/2 A) = 5.7 K only!$

- So we lost a lot of margin from 9.2 K \rightarrow 7.2 K \rightarrow 5.7 K versus 4.4 K.
- At 4.4 K, at 50% I_c and 5 T there is only 1.2 K margin !
- At 75% of I_c we get 0.7 K, so we never can operate very near to I_c !
- Following $\Delta T = Q / c(T)$,

release of energy (heat) from various sources will cause a temperature rise and thus the superconducting state is very seriously in danger.

• The heat that can be absorbed without reaching T_{cs} is the enthalpy difference $\Delta H = \int c(T) dT$ between T_{cs} and T_{o} .

Adiabatic filament stability, d_{fil}

Field penetration in filaments, the Critical State Model

- In the filament magnetic energy is stored
- When disturbed, the heat must be taken up by the enthalpy of the filament
- A disturbance $\Delta T1$ will cause a $-\Delta Jc$, so flux motion, leading to E, this leading to heat and so again a $\Delta T2$
- When ∆T2 > ∆T1, the process will accelerate and the flux profile collapses
- Based on simple slab model, the adiabatic stability criterion is found:

 d_{fil} . $J_c < (3 c (T_c - T_o) / \mu_o)^{1/2}$

So we see a maximum filament thickness for a given current density, to guarantee stability.

For NbTi, c=5600 J/m³; T_c(5 T)=7.2 K, T_o= 4.2 K and J_c = 3000 A/mm², we find $d_{fil} < 70 \mu m$.







Adiabatic Wire Self field Stability, Dwire

Filaments coupled by self field

- Adiabatic filament stability requires fine filaments in a matrix
- These can be de-coupled for transverse fields by twisting
- But are still fully coupled by the self-field
- Again following the CSM, we see the field penetration profile disturbed by a ∆T
- Field profile has to change, penetrates deeper, causing heat dissipation taken up by the enthalpy up to a certain limit
- Assuming η=sc/total ratio and current density ηJ
- We find for the adiabatic self-field criterion:

 $D_{wire} \eta J < (4c (T_c - T_o)/\mu_o)^{1/2} f (I/Ic)$

where f (I/Ic) = 1/(-0.5 ln(I) - $3/8 + i^2/6 - i^4/8$)

 So we see a maximum wire diameter for a given Jc and I/I Commonly is used 0.7< D_{wire} <1.3 mm in cables.



Self-field Stability: cable examples

ITER cable for central solenoid

- 65 kA at 13.5 T, ~1152 Nb3Sn wires parallel in a twisted multi-stage cable.
- Cable layout with 5 stages: 1x3x4x4x4x6.
- Wire 0.81 mm, filaments 4 µm.
- the strands take all positions in the cable to guarantee equal current sharing.

LHC type Nb3Sn Rutherford cable

- 33 stands single stage twisted.
- 13 kA at 11 T.

ATLAS cable

- Al stabilized 40 strands Rutherford cable.
- 65 kA at 5 T.



~1152 wires ITER Nb3Sn cable



33 wires LHC-type Nb3Sn cable



40 strands ATLAS BT cable

Temperature jumps, low heat capacity

Why is release of heat so critical at 4 K ?

- Heat capacity is strongly T-dependent
- Copper-NbTi composite: Cp(T)= η((6.8/η+43.8)T³+(97.4+69.8 B)T) μJ/mm³K, at 5 T and 40% NbTi in a Cu matrix:
- 2.5 μJ/mm³K at 4.2 K and
- 0.5 μJ/mm³K at 1.9 K !
- 2.5 μJ/mm corresponds to a movement in a 1 mm wire at 5 T, 500 A of 1 μm only!



Heat release of μ J/mm³ has to be avoided, otherwise magnet will quench

- avoid friction and slip-stick by introducing low friction sliding (kapton films wrapped around wires and cables)
- avoid any displacement, vacuum impregnation of coils
- avoid resin cracks, avoid local stress concentrations at bonded surfaces

Point disturbance, MPZ

Minimum Propagation Zone (1-d case)

- How large must the distortion be to get a quench ?
- Consider a wire with current I, heat removal Q along the wire and central zone in normal state (simple, one dimensional case)



Look for length L where heat produced is equal to heat removed:

$$ρ$$
 J² A L ≈ 2 λ (T_c-T_{bath}) / L

$$L = (2\lambda(T_c-T_{bath})/\rho J^2)^{1/2} = MPZ$$

Propagation occurs when L > MPZ and recovery when L < MPZ

Minimum Propagation Zone, MPZ

Examples of MPZ in a various wires

 In a bare NbTi wire or filament: take 5 T; 3000 A/mm²; ρ= 6x10⁻⁷ Ωm; λ= 0.1 W/mK; T_c= 7 K and we find 0.3 μm only, pure NbTi can not be used!



- NbTi with CuNi matrix would give $3 \mu m$ and 0.1 μJ !
- Such wire is extremely sensitive to any heat pulse
- Remedy: reduce ρ by using copper matrix (3x10⁻¹⁰ Ω m, factor 2000 !) and increase λ by using copper (>200 W/mK, factor 2000 again !)

We see how wonderful copper (or AI) is, without copper no sc magnets !

- \checkmark factor 2000 improvement, from μm to few mm and μJ range
- ✓ for a typical LHC cable we get about 15 mm
- and in the ATLAS conductor (600 mm² pure Al and 20 kA) we get about 500 mm !



30

Why magnets need High Current & Cables

Magnetic field and stored energy

 $B \propto N.I$ $E \propto B^2.Volume$ Inductance $L \propto N^2$

- Need safe survival from a quench
- Energy dump within short time before conductor burns out
- \rightarrow Thus low N, high current I

Also $I_{safe} \propto J.E/V_d$, kV-range for V_d ,

with usual current densities this leads to 10-100 kA

Given common strand currents of 100 to 500 A, we need for large scale magnets multi-strand cables with 20-1000 strands! No escape!









0.0001 m³ HF insert 200 A

2 m³ MRI magnet **200-800 A** @ 1-3 T, ~10 MJ

25 m³ ATLAS solenoid 8 kA @ 2T, 40 MJ





1000 m³ ITER magnets **40-70 kA** @ 10-13T, 50 GJ

Request for: High current conductors

200 A HTS tape?

200000000

Single: No! Cabled: may be, but to be developed

65000 A@5T Al-NbTi/Cu?



Yes!



One cannot build large scale magnets from single NbTi-Nb₃Sn-B2212-Y123 wires or tapes.

We need superconductors that can be cabled and survive a quench!

Novel Detector Magnet Superconductors

For the next generation detector magnets, conductors are further developed and reinforced, more stored energy, larger size.



Reinforcing AI-stabilized conductors

- Option 1 Ni or Zn - doped Aluminum
- Used in the ATLAS Solenoid mechanical reinforcement while keeping quench stability





Option 2

Reinforce with Al-alloy side bars, EB-welded to the Al and NbTi/Cu co-extruded conductor

Doable but expensive





Alternative: use a Cable-in-Conduit

More than 25 years cable-in-conduit conductors (CICC) are in use for fusion type of magnets with forced flow helium to maximize heat removal and stability.



Very flexible in choosing cable size, current rating, strength and helium cooling directly on the superconductor -> maximum stability



The energy stored in a magnet is $W_L = \frac{1}{2} L I^2 [J] = \frac{1}{2} \int BH dV$, the energy density being $\frac{1}{2} BH \text{ or } B^2/2\mu_0$

This energy could be absorbed by the magnet cold mass assuming a safe temperature T_m

- $W_L/m =_o \int^{Tm} C_p(T) dT = H(T_m) H(T_o = 4.2)$ $\approx H(T_m)$ since $C_p(4.2)$ is negligible
- For 150 K, we can absorb about 20 kJ/kg cold mass provided uniformly distributed
- Usual values for W_L/m are in the range
 <10 kJ/kg, so apparently no problem
- But heat distribution must be controlling the normal zone spatial distribution and speed.


Adiabatic heating of the conductor

Temperature of the conductor?

 Heating in the normal zone ρJ² is taken up by the conductor enthalpy:

 $\rho(T) J^2(t) dt = c(T) dT$

 $_{o}\int ^{t}J^{2}(t) dt =_{4}\int ^{T}c(T)/\rho(T)dT = constant = F(T_{m})$

- F is the Load Integral, used to assess transient thermal loads in devices.
- F is a constant, calculated for NbTi, Cu, resin and any mixture as a winding.
- Typical values for F(T_m) are in the range 2-9x10¹⁶ for 150 K and 5-15 for 300 K maximum temperature depending on the conductor composition.





TEMPERATURE 0 (K)

Adiabatic hot spot temperature

 $_{o}\int^{t} J^{2}(t) dt =_{4}\int^{T} c(T)/\rho(T) dT = constant = F(T_{m})$

Simple solutions exist for constant or exponential decaying currents

Constant current

 $J^2 t_m = F(T_m) \rightarrow t_m < F/J^2$

Exponential decay

 $J^2 \tau / 2 = F(T_m) \rightarrow \tau < 2F/J^2$

Examples

- NbTi/Cu and CuNi matrix conductors with J = 500 A/mm²
- F(300) ∞ 1/ ρ
- F(300) for Cu is ~1.4 10¹⁷ and ~1.4 10¹⁶ for CuNi (or pure NbTi)
- Maximum τ in NbTi/Cu before reaching 300 K is a 0.1-1 second
- Maximum τ in NbTi or NbTi/CuNi is ~ms, so very little time to react and the conductor will burn out when used at high current density !



Safe hot spot temperature

Criterion for hot spot temperature

- Beyond 900 K AI structures start to collapse.
- Beyond 650 K we start to lose pinning, so J_c .
- Even 300 K is too high, as it endangers the windings.
- Severe thermal shock due to differential thermal contractions will occur.
- This may cause resin cracking and debonding, and thus training or degradation.
- A "safe" hot spot temperature is 100-150 K!
- Usually 100 K is taken nominally and a peak of 200-300 K for exceptional cases (failing protection systems for example).
- ✤ 300 K may be acceptable for an R&D magnet, but is not an acceptable design value for a detector magnet that has to survive, operate at minimum risk and must be quench-recovered within 3-4 days.



1.2

1.0

0.8

0.6

0.4

٥

•/。

CONTRACTION

Destructive power of uncontrolled quenches

LHC dipole of 15m and 8.35T stores 8 MJ, which corresponds to melting 1.5L of copper, enough to evaporate 10cm of coil !

And we have seen in Sep 2008 what a few magnet quenches can do!

ATLAS detector toroid stores 1.6 GJ, good for 600L of melted copper, or equivalent to the collision energy of 100 trucks of 40 tons with speed of 100 km/h!

To be safe with equipment and personnel: Quench Protection has to cover all possible quenches in the entire electrical circuit from + to – terminal on the cryostat (current leads & bus connections & coil).



Damage at an LHC interconnect





Quench detection circuit

- The magnet safety system comprises the quench detectors, logics for opening switches and to supply current to the quench heaters.
- The system must be extremely reliable and power secured.
- The motto is : "keep it simple", meaning robust and straight forward detection circuits, simple electronics, hardwired and 3-5 times redundant.
- First the quench, a normal zone, must be detected, then switches have to be opened and quench heaters activated.



Quench detection methods

Bridge method

- Detects the resistance in any branch of the coils, very robust, simple and proven.
- 3 sets of bridges, asymmetrically connected to see symmetric quenches.
- Commonly used for large magnets.

Voltage across coil

 Voltage across coil compensated for the inductive component. Requires differential amplifiers, more complicated, more electronics.

Other methods

- Temperature, pressure gages, pick-up coils, strain sensor, etc.
- Many proposed, but mostly not used.





Toroids quench detection

- 1.5 GJ energy, 20 kA current, 4 T peak field, 3 kJ/kg stored
- 3 toroids, each comprising 8 flat coils, thermally not connected
- 22 m diameter
- 5 m x 26 m long coils



Detector characteristics



- All toroids 3 x 8 = 24 coils are connected in series.
- The energy is dumped in the 3 toroid cold masses, voltage limited to 40V.
- Quench detection by 3 bridges + 3 differential units per toroid so 6 fold redundancy, heaters are fired introducing 4 normal zones in every coil, expected maximum hot spot temperature ~100K.
- Threshold 0.3 V
- Low pass filter 1 s
- Fast dump in about 80 s.







Toroid Fast Dump test result:

- Provoked Quenches at 20.5 kA, heaters fired, quench is spread
- ~ 60 K cold mass temperature at 20.5 kA, recovery in about 80 hours
- ~ 90 K hot spot in the conductor, perfectly safe quench behavior.





3. Designing a detector magnet, example CMS solenoid



Design steps: example CMS solenoid

- 1. Magnetic field calculation
- 2. Effect of the iron yoke
- 3. Magnetic stored energy
- 4. Lorentz forces in the coils
- 5. Hoop stress
- 6. Choosing current vs selfinductance
- 7. Conductor dimensions and layers
- 8. Conductor details
- 9. Stabilizer, Cu or Al





Design steps: Magnetic field, no iron

Field calculation without iron yoke:			
Current density: $J = \frac{NI}{L(b-a)}$			
Field $B_o = Jr\mu_o\beta\left\{\frac{\alpha+\sqrt{(\alpha^2+\beta^2)}}{1+\sqrt{1+\beta^2)}}\right\}$			
$B_o = \mu_o n I$ for $\beta \to \infty$			

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With real CMS magnet sizes: r = 3200 mm; R = 3418 mm L = 12500 mmN = 2180; I = 19500 A



We find:
$$B_o(\alpha, \beta) = 3.77 T$$
 (88% of infinite)
 $B_o(\beta = \infty) = 4.27 T$

With a FEM code we find 3.77 T as well



3.772e+000 : >3.970e+000

Design steps: Magnetic field, with iron

Accurate analytical formulae do not exist, a calculation with a FEM code is needed (OPERA-3D, ANSYS, COMSOL).

Simple solid magnetic yoke:

 $B_o = 4.17 T$ (98% of infinite)



Iron is a magnetic mirror, the coil is almost infinite.

Real iron with gaps for detectors:

 $B_o = 4.0 T$ in center

4.6 T in conductor

Stored energy:

FEM calculation yields: $\frac{1}{2\mu_o} \int B^2(r,z) dV = 2.6 GJ$ Simple approximation: $\frac{1}{2\mu_o} B^2 V = 2.46 \text{ GJ}$, V = bore volume

4.104e+000 : >4.320e+000
3.888e+000:4.104e+000
3.672e+000 : 3.888e+000
3.456e+000 : 3.672e+000
3.240e+000 : 3.456e+000
3.024e+000 : 3.240e+000
2.808e+000:3.024e+000
2.592e+000 : 2.808e+000
2.376e+000 : 2.592e+000
2.160e+000 : 2.376e+000
1.944e+000 : 2.160e+000
1.728e+000 : 1.944e+000
1.512e+000 : 1.728e+000
1.296e+000 : 1.512e+000
1.080e+000 : 1.296e+000
8.642e-001 : 1.080e+000
6.481e-001 : 8.642e-001
4.321e-001 : 6.481e-001
2.161e-001 : 4.321e-001
<7.652e-005 : 2.161e-001

4.6859+000 : >4.932e+000 4.4389+000 : 4.6859+000 4.1929+000 : 4.6859+000 3.9459+000 : 4.1929+000 3.6999+000 : 3.9459+000 3.4529+000 : 3.6999+000 2.9599+000 : 3.2069+000 2.7129+000 : 3.2069+000 2.4669+000 : 2.7129+000 2.2199+000 : 2.2199+000 1.9739+000 : 2.2199+000 1.9739+000 : 2.2199+000 1.4799+000 : 1.726e+000 1.233e+000 : 1.479e+000 9.863e-001 : 1.233e+000 7.397e-001 : 9.863e-001 4.932e-001 : 7.397e-001 2.4669-001 : 4.932e-001 2.4669-001 : 4.932e-001		
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3.945e+000 : 4.192e+000 3.699e+000 : 3.945e+000 3.452e+000 : 3.699e+000 3.206e+000 : 3.452e+000 2.959e+000 : 3.206e+000 2.712e+000 : 2.959e+000 2.466e+000 : 2.712e+000 2.219e+000 : 2.466e+000 1.973e+000 : 2.219e+000 1.726e+000 : 1.973e+000 1.479e+000 : 1.726e+000 1.23e+000 : 1.479e+000 9.863e-001 : 1.23e+000 7.397e-001 : 9.863e-001 4.932e-001 : 7.397e-001 2.466e-001 : 4.932e-001		4.438e+000 : 4.685e+000
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Design steps: Magnetic forces

Lorentz forces due to B and J cause axial compressive forces and radial forces causing hoop stress:

 $\overline{F} = \int (\overline{J} x \overline{B}) dV$

- Radial field causes axial force F_a
- Axial field causes radial forces F_r
- In fact the solenoid wants to blow up into a ball shape

For CMS: $F_a = +1.66 \text{ GN}$, $F_r = -140 \text{ MN} (14 \text{ kt})$

The "Ball" **Pressure** \approx **F**_r/surface = 6.6 MPa

- Magnetic pressure = ${}^{B^2}/_{2\mu_o} = 6.4 MPa$
- or 64 atm



Design steps: Hoop stress, coil thickness

The radial pressure is reacted in the cylinder with thickness t (windings + extra material) by the hoop stress:

$$\sigma_{hoop} = {^{a P_r}/_t}$$

To be respected design rule:

 $\sigma_{hoop,max} = 2/3 \rho_{yield}$

Structural coil thickness:

$$t = \frac{3 r P_r}{2 \rho_{yield}} = 320 \ mm$$
 ,

using 100 MPa annealed Al5083, or

- t = 190 mm, based on special 170 MPa Al5083-H321.
- So we need some 190 320 mm thick structural special AI alloy on top of the soft conductor to withstand the radial forces in a safe way.



Design steps: Current vs self-inductance

Self-inductance L_c and current I are linked through the stored energy:

$$E = \frac{L_c I^2}{2} = \frac{1}{2\mu_o} \int B^2 \, dV \approx \frac{1}{2\mu_o} B_o^2 V, \text{ and } L_c = \mu_o N^2 \pi r^2 2/L$$

- Current I must be high for protection reasons, say 20 kA
- Then $L_c \approx 14$ H and for N follows N ≈ 2100.
- Adaptation to conductor & coil dimensions leads to 19.5 kA / 2180 turns.
- The coil has 42.5 10⁶ ampere-turns.

In the windings section of

- ≈ 320 mm x 12500 mm we have to put in place:
- 2180 turns of superconducting cable with 19.5 kA
- o extra stabilizing and quench protection material around the cable
- conductor insulation
- structural reinforcement for handling the hoop stress
- an outer support cylinder for integrity and conduction cooling supply



53

Design steps: Conductor size and layers

- 4 T is made with 2180 turns and 19.5 kA current, but: How many layers is wise?
- Coil winding section is 12500 mm x 263 mm,
- n layers x conductor height = 263 mm
- Use 1 (easy), or even number of layers: 2, 4 or 6
- 1 or 2 layers requires a too thin conductor to be wound on its small edge.
- Then 4 layers is best, few layers only and acceptable conductor size of 66 x 23 mm², 6 layers would mean 44 x 34, almost square.

There is a thermal argument as well:

 winding on small-edge gives less layers, so less thick insulation (resin, glass, polyimide) between the superconductor (NbTi) and the heat sink (cooling pipe), thus a small temperature gradient.





Design steps: Superconductor needed

The coil runs at 19.5 kA with a peak field of 4.6 T at 4.5 K:

- Critical current density at 4.6 T/4.5K including 5% cabling degradation is 3000 A/mm².
- We need margin so we run at 1/3 of the critical current, at 1000 A/mm².



- 19500 A and 1000 A/mm², \rightarrow need 19.5 A/mm² sc per turn=cable
- Self-field stability \rightarrow wire diameter <1.28 mm
- A minimum Cu/sc ratio is $1:1/1 \rightarrow Asc= 0.61 \text{ mm}^2$
- Number of strands in the cable is then 19.5/0.61 = 32.
- Filament size? Adiabatic filament stability requires <40um.
- The filament section is 0.00126 mm² \rightarrow we need \geq 484 filaments.
- Twist pitches on strand a cables can be standard giving a good cable stability as needed for the cable/Al co-extrusion process.
- Thus Ls=25 mm and Lc= 185 mm and twist directions SZ.

Design steps: wire & cable specification

Following these arguments the cable specification is now as follows:







Strand Constituents	Material
High homogeneity Nb-Ti	Nb 47±1 W t % Ti
High Purity Copper	RRR > 300
Niobium Barrier	Reactor Grade I
Strand Design Parameters	Parameters
Strand Diameter	$1.280\pm0.005~\mathrm{mm}$
(Cu+Barrier)/Nb-Ti ratio	1.1 ± 0.1
Filament diameter (mm)	< 40
Number of Filaments	• 552
Strand Unit length (m)	2750
Twist Pitch	$45 \pm 5 \text{ mm } Z \text{ (RHS)}$
Strand Minimum Critical Current Ic (A)	1925
(Criteria : 5 T, 4.2 K, 10 µV/m)	
<i>n</i> -value 5T	>40
Final copper RRR	>100
Rutherford cable	
Cabling direction	S
Nominal current	19500 A
Critical current at 5T, 4.2K	≥56000 A
Critical temperature at 4.6T	7.35 K
Current sharing temperature at 4.6T and 19.5 kA	≥6.33 K
strand number	32
dimensions	20.68x2.34 mm ²
Cable transposition pitch	185 mm
Cable compacting ratio	87 %

Design steps: Cable - Al co-extrusion

The cable is co-extruded with high purity AI (RRR>1500)



Coil windings build up

Now we have: 4 layers of a soft conductor Al/NbTi/Cu, 127 mm thick and a thick support cylinder of 186 mm.

Is this thermally and mechanically an optimal design? No !









Making of CMS Solenoid: support cylinder

The CMS magnet cold mass was made in 5 units mostly at ASG – Genua, transported to CERN for on-surface assembly and then insertion as a whole in the CMS cavern.



Support cylinder manufacturing, 5 units



Thermal siphon cooling layout, pipework welded to the cylinder

Making of CMS Solenoid: coil winding



Bend conductor pressed against cylinder



Conductor spiral leading into cylinder

Dedicated coil winding machine allowing winding inside the support cylinder (6.2 m diameter)





Conductor bending



Taping insulation on conductor

Making of CMS Solenoid: vac impregnation









Vacuum impregnation tools, resin curing, result: Clear transparent resin





Making of CMS Solenoid: assembly on site





Modules transport, stacking, integration in cryostat and finished coil ready for insertion in cavern. READY !



4. The making of ATLAS.....



ATLAS on surface and underground



height = 35 m

- Underground cavern at - 90 m
 - 2 shafts give access to a 50,000 m³ cavern for the detector

64

ATLAS at the White House





ATLAS sc magnet system

- 1 Barrel Toroid, 2 End Cap Toroids and 1 Central Solenoid
- 4 magnets provide 2 T magnetic field for the inner detector (solenoid) and ~1 T for the muon detectors in blue (toroids)
- 20 m diameter x 25 m long
- 8300 m³ volume with field
- 170 t superconductor
- 700 t cold mass
- 1320 t magnets
- 7000 t detector
- 90 km superconductor
- 20.5 kA at 4.1 T
- 1.6 GJ stored energy
- 4.7 K conduction cooled
- 9 yrs of construction 98-07





Magnetic field configuration

 2 T in Solenoid closed via return yoke 2.6 T peak in windings
~ 0.8 T average in Barrel Toroid torus 3.9 T peak in windings
~ 1.3 T average in End Cap Toroid 4.1 T peak in windings











ATLAS Barrel Toroid Integration

Construction of a single coil, 8 of these constitute the toroid

Two racetrack double pancakes



- 2 x 60 turns, pre-stressed and glued in an AI 5083 casing
- Forced flow indirect cooling via redundant circuits of AI 1050 alloy tubes glued on the casing
- Al alloy thermal shield panels
- Superinsulation
- 8 Ti Tie rods
- 16 fre lateral supports
- Instrumentation
- SS vacuum vessel
- Al-alloy warm structure



ATLAS: manufacturing the parts



ATLAS: Toroid coils integration and test



ATLAS: Start of Barrel Toroid assembly







 Transport, decent, reception
Complex but safe manipulations
Lowering using 2 lifting frames
Hydraulic winch with load capacity 190 t (subcontracted)
ATLAS: BT method of installation



Assemble in egg shape with dy=+30mm.

Coils put in calculated positions.

Keep all coils in position and fill the gaps between coils.

When toroid is closed, take away supports.

Put all other mass on.

Finally shape will become cylindrical dy~0 mm

ATLAS: Barrel Toroid assembly coils 1-3



✓ First the 2 coils in the feet then the other 6

✓ and a lot of temporary (green) support structures to position the coils in space



ATLAS: Barrel Toroid assembly coils 4-8



ATLAS: Barrel Toroid in cavern (Nov 05)



CERN

ATLAS: Two End-Cap Toroids

Aret Toroid Ner Detector Ladorimeter 2 x 8 coils

End Cap To

 $4 \times 4.5 \text{ m}^2$

20 kA, 4.1 T peak

Al 5083 cold mass, torus assembly, 8 keystone boxes hanging on bore tube

Al 5083 vacuum vessel

Size 11m dia x 5m





ATLAS: End Cap Toroids on the move.....



Magnet system services: isolation vacuum



- 4 backing pumps, 21 diffusion pumps, stops when no water cooling and power
- must run 24/7, on UPS & diesel, redundant water cooling circuits

Magnet services: helium cryogenics



ATLAS: He proximity cryogenics



Magnet services: current, 20.4 kA – 18 V



Toroids in series:

- dump in parallel
- power convertor
- 2 switches
- dump resistors
- diode units
- 240 m Al bus bars







Barrel Toroid test in Nov 2006

- Few months of chasing leaks, repairs & vacuum cleaning
- Cooling down (340 t) took 5 wks with shield refrigerator to 70K, then with main refrigerator to 4.6K
- Helium circulation pumps for 10L/s, 1200g/s and work great
- ✓ No surprises in coil mechanics
- ✓ Test: in steps to 20.5kA nominal, to 21kA to prove margin, provoke heater induced quench
 → fast dump....
- ✓ Tmax cold mass = 58K
- Thot spot = ~85K, very safe !
- Barrel Toroid accepted



ATLAS: From Concept to Commissioning

- Concepts, seeking consensus, predesign
- Construction approval
- Industrial components production
- Integration, on surface test & installation
- Test and commissioning
- Stable operation, ready for physics
- 1st repair of LHC after splice incident
- First 3yrs physics data taking period
- First long shut down, consolidation works
- And another15-20 yrs depending on physics results......
- In total 17 years from predesign to ready for physics
- And expected operational life time of ~25 years

1991 - 1994 Sep 1997 1998 - 2005 2002 - 2006 2007 - 2008 Aug 2008 Sep 08 - Aug 09 Sep 09 - Feb 13 Mar 13 - Jan 15





	Т	otal in MCHF	
	Barrel Toroid	80	
	End Cap Toroids	37	
	Central Solenoid	11	
	Vacuum, Cryogenics, Current & Controls	<u>31</u> +	
	Recognized total cost by ATLAS	159 MCHF	
	Initial budget, no reserve, no inflation correction	137	
✓	Extra cost across 10 yrs of construction, only:	22 (16%)	
\checkmark	~ 65% was financed and produced as in-kind contributions, worked fine!		
	Free contributions, hidden manpower,		
	and cost savings through simplifications:	~ 40	
	True cost of original design (already anticipated in 1996!):	~ 200 MCHF	
\checkmark	Financially the project was concluded in a satisfactory way	,	

Operation Statistics since Sep 2009

Magnet	Ramps	Slow Dump	Fast Dump	Quench
Solenoid	57	54	4	0
Toroids	74	69	5	4
	Solenoid		Toroids	
Effective field-ON time	tive field-ON time 834 days		829 days	
Percentage of time ON (since data taking Sep 09)	67 %		67 %	

- Magnet services, pumps, cryogenics, controls run since Jan 2006
- Magnets commissioned in August 2008
- In operation for collisions since September 2009 \rightarrow ~2040 ?
- In first 2 years many stops for adapting magnet services or to do detectors repairs, will improve
- 5 Fast Dumps in Toroid and 3 in Solenoid
- So far no quenches originated in coils, in current leads only











WS POLITIEK OPINIE BUITENLAND SPORT TECH & MEDIA

Volkskrantnl

July 4, 2012 **CERN** press conference



পেয়েছি, যা খঁজছিলাম





VOLUME 13, NUMBER 16

PHYSICAL REVIEW LETTERS

19 October 1964

BROKEN SYMMETRIES AND THE MASSES OF GAUGE BOSONS

Peter W. Higgs Tait Institute of Mathematical Physics, University of Edinburgh, Edinburgh, Scotland (Received 31 August 1964)





Physics Letters B Volume 716, Issue 1, 17 September 2012, Pages 1–29



Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC $\dot{\star}$

Universally Available



"I certainly had no idea it would happen in my lifetime at the beginning, more than 40 years ago. I think it shows amazing dedication by the young people involved with these colossal collaborations to persist in this way, on what is a really a very difficult task. I congratulate them."

Peter Higgs, July 4th, 2012



5. Detector Magnets for a 100 TeV p-p collider

Future Circular Collider study Design drivers Three options considered Twin Solenoid features Conclusion



Options for increasing colliding energy

Energy = 0.3 x B x R

B: 1.8 x from NbTi to Nb₃Sn
B: 2.4 x from NbTi to HTS
R: 4-5 x more magnets

- New 80-100 km tunnel in Geneva area
- pp-collider (VHE-LHC) defining the size
- Options for adding an e+e collider (TLEP) p-e collider (VLHeC)
- CERN-hosted study with international collaboration



It easily takes 30 years time..... start now

"CERN should undertake design studies for accelerator projects in a global context, with emphasis on **proton-proton** and electron- positron **highenergy frontier machines.**"



FCC Study : p-p towards 100 TeV

Kick-off meeting already happened, mid-February 2014

Design drivers for detector magnets

Bending power: 100 TeV, a 7 x higher collision energy than 14

Same tracking resolution

BL² has to be increased by factor 7!

- $\frac{\sigma(p_T)}{p_T} = \frac{\sigma(\kappa)}{\kappa} = \frac{\sigma_x \cdot p_T}{0.3BL^2} \sqrt{\frac{720}{(N+4)}}$
- In single solenoid: increase field up to 6 T
 In solenoid-toroid system: in solenoid around the ID, need a field/track length combination of 3.5T/3m or 2T/4m, and a toroid with ≈2 T and 1.5 x increase of tracking length.

Also need low-angle coverage in forward direction

> add a dipole or iron toroid for on-beam bending featuring some 10 Tm!

HCAL depth increase from 10 λ to 12 λ (iron) radial thickness some 3.0 m!

 \succ Free bore of solenoid or toroid increases to 6 m and length accordingly.

ECAL to cover low angles, move out, from 5 to 15 m, system gets longer.

Higher magnetic field, larger bore, longer system. 3 options studied.

Option 1: Solenoid – Yoke + Dipoles (CMS inspired)



Solenoid: 10-12 m diameter, 5-6 T, 23 m long
 + massive Iron yoke for flux shielding and muon tagging.

Dipoles: 10 Tm with return yoke placed at z≈18 m. Practically no coupling between dipoles and solenoid. They can be designed independently at first.

Option 1: Solenoid in Yoke + Dipoles



6 T in a 12 m bore, 23 m long, 28 m outer diameter.

Stored energy 54 GJ, 6.3 T peak field.

Yoke: 6.3 m thick iron needed to have 10 mT line at 22 m, 15 m³, mass ≈ 120,000 ton !!! (>300 M€ raw material).

Huge mass, serious consequences for cavern floor, installation, opening -closing system, bulky, not an elegant design.

Option 2: Twin Solenoid + Dipoles



Twin Solenoid: a 6 T, 12 m dia x 23 m long main solenoid + an active shielding coil

Important advantages:

✓ **Nice Muon tracking space:** area with 2 - 3 T for tracking in 4-5 layers.

✓ Very light: 2 coils + structures, \approx 5 kt, only \approx 4% of the option with yoke!

✓ Much smaller: system outer diameter is significantly less than with iron.



Example:

Main solenoid:

- 6.0 T in 12 m bore, 12 m long,
- 6.3 T peak field, 10 A/mm²

Shielding solenoid:

≈ 3 T in 3.5 m gap

22 m bore, 28 m long, 10 A/mm 2

Mass:

≈ 2 kt main coil + ≈1.8 kt shield coil

in total with supports some 4-5 kt only!

- Nice gap for muon tracking: 3.5 m gap with 3 T (local ≈10 Tm or ≈35 Tm²).
- Shielding: 5 mT line at 34 m from center.





Option 3: Toroids + Solenoid + Dipoles (ATLAS +)





- 1 Air core Barrel Toroid with 7 x muon bending power $B_z L^2$.
- 2 End Cap Toroids to cover medium angle forward direction.
- 2 Dipoles to cover low-angle forward direction.
- \diamond Overall dimensions: 30 m diameter x 51 m length (36,000 m³).

Option 3: Toroids + Solenoid + Dipoles



- 10 coils in Barrel Toroid and 2 x 10 coils in End Cap Toroids.
- Peak field on the conductor ≈6.5 T for 16 Tm and ≈8 T for 20 Tm, to be minimized by locally reshaping the coil or reduction of current density.
- Can still be done with NbTi technology (to limit cost)!

Option 3: Toroids + Solenoid + Dipoles



- 3.5 T in central solenoid, 2 T 10 Tm in dipoles and ≈1.7 T in toroid.
- ✤ 55 GJ stored energy (for 16 Tm; 130 Tm²)!
- 0.6 GJ in Solenoid, 0.9 GJ in 2 Dipoles, 2x2.1 GJ in the two End Cap Toroids, and 47.5 GJ in the Barrel Toroid.

Superconductors - change of technology

Peak magnetic field of 7 to 8 T implies high winding stress and a low temperature margin, just in reach of NhTi provided correctly cooled.

just in reach of NbTi provided correctly cooled.

- Classical Ni doped Al-stabilized NbTi Rutherford cable may be used for the "small" 3.5 T, 4 m bore solenoid requiring high transparency.
- All other coils require highstrength materials and direct cooling of the superconductor, asking for use of cable-inconduit type of conductor.









Sizes: 12 m bore, 30 m dia, 30 to 50 m length.

- Looks gigantic but similar sized magnets are being made these days (ITER PF coils, 26m).
- Production on site, in smaller modules, but very well possible.

Stored Energy: in 50 to 100 GJ range

- High values but doable.
- Combination of energy extraction and dump in cold mass, controlled by a redundant, fail-safe quench protection system.

There are no principle technical problems impeding the constructing of these magnets.









This concludes the course Enjoy the rest of the day.....

Contraction of the second

Presented: 1. Concepts

- Superconductors
 Design of the CMS solenoid
 The making of ATLAS
 - 5. Future Collider Detectors

